

Relativistic Corrections to the Sunyaev-Zeldovich Effect for Clusters of Galaxies

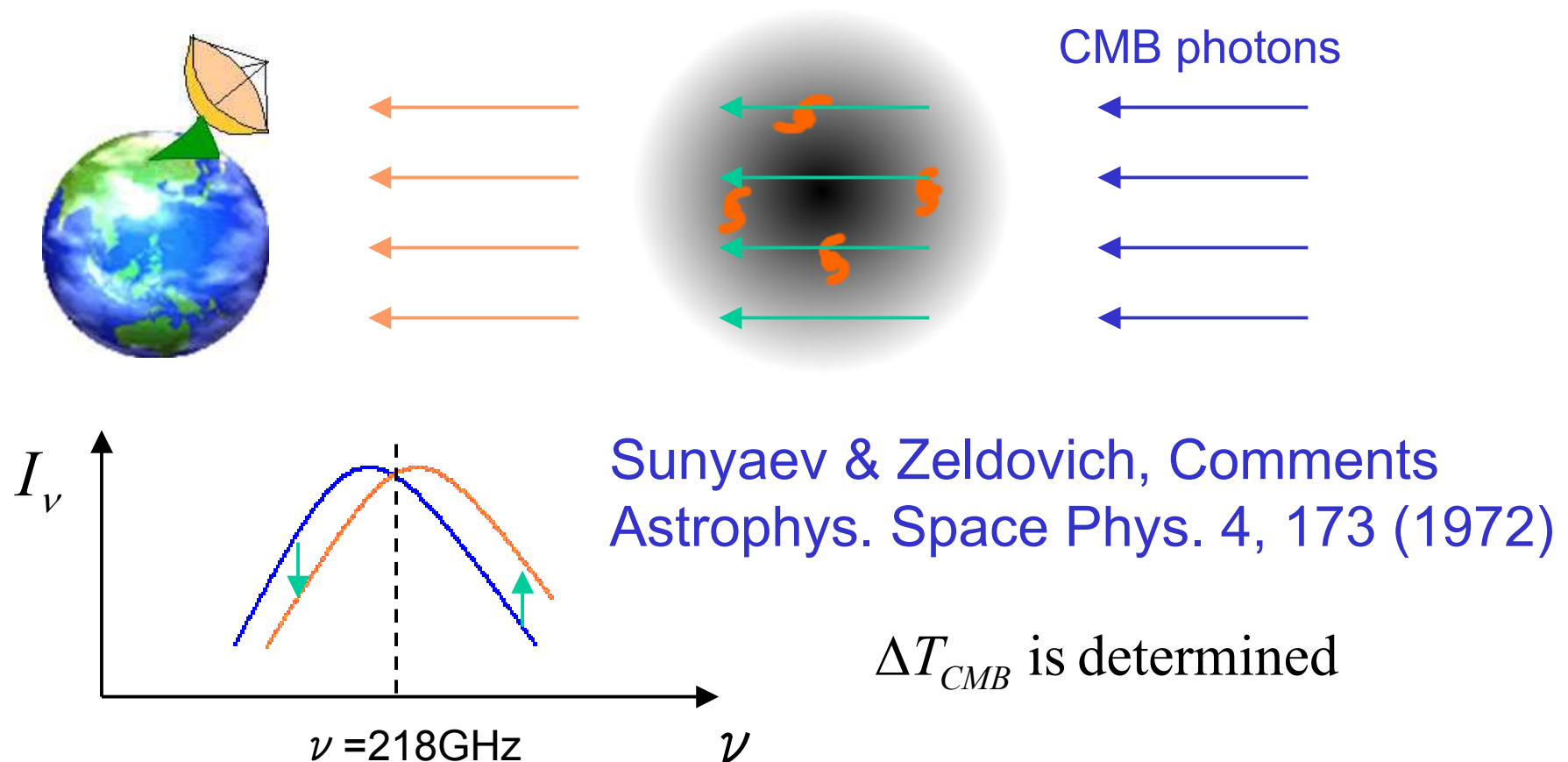
Satoshi Nozawa
Josai Junior College for Women
Collaborators: N. Itoh, Y.
Kohyama

Contents

- (1) Introduction
- (2) Covariant Formalism for SZ Effects
 - (a) Thermal SZ Effect
 - (b) Kinematical SZ Effect
 - Motion of Clusters of Galaxies
 - Motion of Observer (Solar System)
 - (c) Polarization SZ Effect
 - (d) Double Scattering Effect
- (3) Summary

(1) Introduction

Distortion of the Cosmic Microwave Background (CMB) photon spectrum due to electrons in the CG.
(Inverse Compton scattering)



Determination of the Hubble constant with CG

$$\begin{cases} \Delta T_{CMB} \propto N_e l T_e & \text{--- Thermal SZ effect} \\ S_{X-ray} \propto N_e^2 l T_e^{1/2} & \text{--- X-ray intensity} \end{cases}$$

$$\therefore l \propto (\Delta T_{CMB})^2 S_{X-ray} T_e^{-3/2}$$

l : diameter of CG along the line of sight

$$D_A = \frac{l}{\theta} \text{ (angular diameter distance)}$$

With the redshift observation, the **Hubble constant** is determined ($H_0 = v/D_A$).

Recent observations:

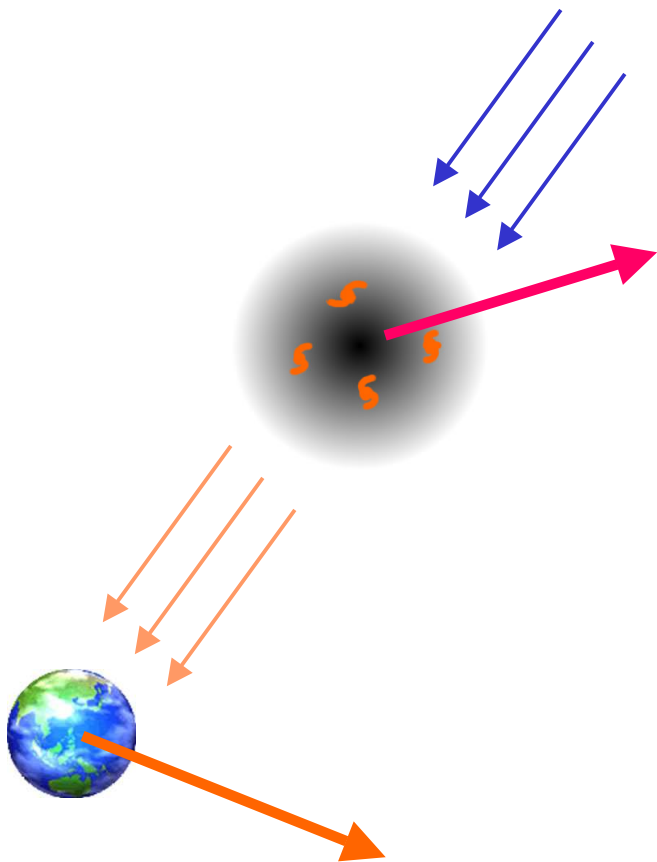
High temperature CG $T_e = 10 \sim 20 \text{keV}$.

Relativistic corrections become significant!

Typical order of magnitude

$$\theta_e \equiv T_e / m_e = 10 / 511 \approx 0.02$$

(2) Covariant Formalism for SZ Effects



- ① Boltzman equation (CMB frame)
→ Thermal SZ
- ② Lorenz boosting CG
peculiar motion of CG
→ Kinematical SZ (CG)
- ③ Lorenz boosting the observer
motion of observer
→ Kinematical SZ (observer)
- ④ polarization cross sections & Stokes parameters for CMB
→ Polarization SZ
- ⑤ Iteration of Boltzman equation
→ Double scattering SZ

(a) Thermal SZ Effect

Itoh, Kohyama & Nozawa, ApJ 502, 7 (1998)

Boltzman equation for the photon distribution function.

$$\frac{\partial n(\omega)}{\partial t} = -2 \int \frac{d^3 p}{(2\pi)^3} d^3 p' d^3 k' W \{n(\omega)[1 + n(\omega')]f(E) - n(\omega')[1 + n(\omega)]f(E')\},$$

$$W = \frac{(e^2/4\pi)^2}{2\omega\omega'EE'} \bar{R} \delta^4(p + k - p' - k'),$$

\bar{R} : Inverse Compton scattering amplitude

$f(E), f(E')$: Fermi distribution functions

p, p' : electron momenta, k, k' : photon momenta

Fokker-Planck
expansion

$(\Delta x \ll 1)$

$$\Delta x \equiv \frac{\omega' - \omega}{k_B T_e}, \quad x \equiv \frac{\omega}{k_B T_e}$$

$$\begin{aligned} \frac{\partial n(\omega)}{\partial t} = & 2 \left[\frac{\partial n}{\partial x} + n(1+n) \right] I_1 \\ & + 2 \left[\frac{\partial^2 n}{\partial x^2} + 2(1+n) \frac{\partial n}{\partial x} + n(1+n) \right] I_2 \\ & + 2 \left[\frac{\partial^3 n}{\partial x^3} + 3(1+n) \frac{\partial^2 n}{\partial x^2} + 3(1+n) \frac{\partial n}{\partial x} + n(1+n) \right] I_3 \\ & + 2 \left[\frac{\partial^4 n}{\partial x^4} + 4(1+n) \frac{\partial^3 n}{\partial x^3} + 6(1+n) \frac{\partial^2 n}{\partial x^2} + 4(1+n) \frac{\partial n}{\partial x} + n(1+n) \right] I_4 \\ & + \dots, \end{aligned}$$

$$I_k \equiv \frac{1}{k!} \int \frac{d^3 p}{(2\pi)^3} d^3 p' d^3 k' W f(E) (\Delta x)^k.$$

Distortion of the CMB spectrum

Expansion in terms of $\theta_e \equiv \frac{k_B T_e}{m_e c^2}$

$$\frac{\Delta n(X)}{n_0(X)} = \frac{\tau \theta_e X e^X}{e^X - 1} [Y_0 + \theta_e Y_1 + \theta_e^2 Y_2 + \theta_e^3 Y_3 + \theta_e^4 Y_4] ,$$

Relativistic corrections

$$X = \frac{\hbar \omega}{k_B T_0} , \quad \tau \equiv \sigma_T \int dl N_e , \quad \sigma_T : \text{Thomson crosssection}$$

T_0 : CMB temperature

$$Y_0 = -4 + \tilde{X} , \quad \text{--- Non-relativistic term}$$

$$\tilde{X} \equiv X \coth\left(\frac{X}{2}\right)$$

Relativistic corrections

$$Y_1 = -10 + \frac{47}{2}\tilde{X} - \frac{42}{5}\tilde{X}^2 + \frac{7}{10}\tilde{X}^3 + \tilde{S}^2\left(-\frac{21}{5} + \frac{7}{5}\tilde{X}\right),$$

$$Y_2 = -\frac{15}{2} + \frac{1023}{8}\tilde{X} - \frac{868}{5}\tilde{X}^2 + \frac{329}{5}\tilde{X}^3 - \frac{44}{5}\tilde{X}^4 + \frac{11}{30}\tilde{X}^5$$

$$+ \tilde{S}^2\left(-\frac{434}{5} + \frac{658}{5}\tilde{X} - \frac{242}{5}\tilde{X}^2 + \frac{143}{30}\tilde{X}^3\right)$$

$$+ \tilde{S}^4\left(-\frac{44}{5} + \frac{187}{60}\tilde{X}\right),$$

$$Y_3 = \frac{15}{2} + \frac{2505}{8}\tilde{X} - \frac{7098}{5}\tilde{X}^2 + \frac{14253}{10}\tilde{X}^3 - \frac{18594}{35}\tilde{X}^4$$

$$+ \frac{12059}{140}\tilde{X}^5 - \frac{128}{21}\tilde{X}^6 + \frac{16}{105}\tilde{X}^7$$

$$+ \tilde{S}^2\left(-\frac{7098}{10} + \frac{14253}{5}\tilde{X} - \frac{102267}{35}\tilde{X}^2 + \frac{156767}{140}\tilde{X}^3 - \frac{1216}{7}\tilde{X}^4 + \frac{64}{7}\tilde{X}^5\right)$$

$$+ \tilde{S}^4\left(-\frac{18594}{35} + \frac{205003}{280}\tilde{X} - \frac{1920}{7}\tilde{X}^2 + \frac{1024}{35}\tilde{X}^3\right)$$

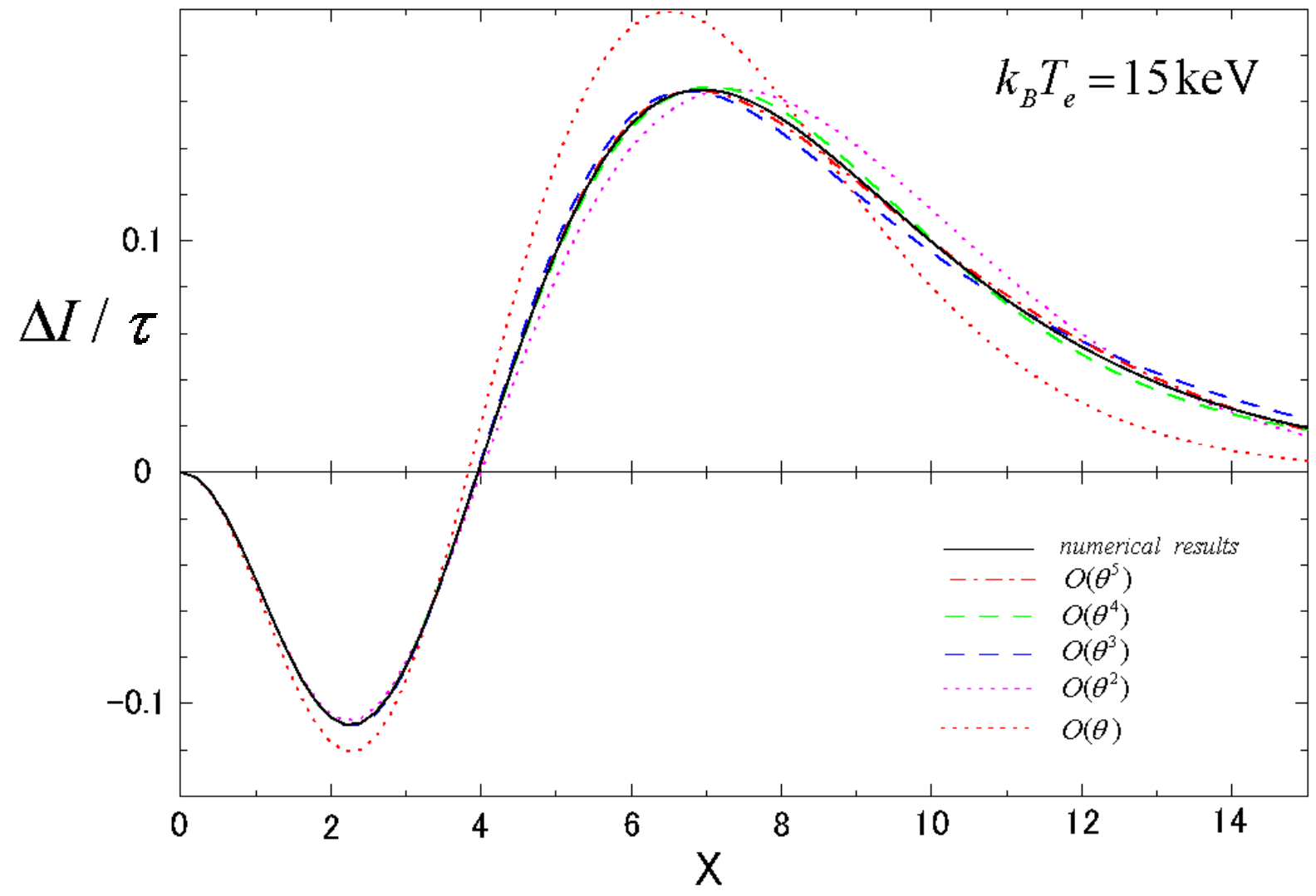
$$+ \tilde{S}^6\left(-\frac{544}{21} + \frac{992}{105}\tilde{X}\right),$$

$$\begin{aligned}
Y_4 = & -\frac{135}{32} + \frac{30375}{128}\tilde{X} - \frac{62391}{10}\tilde{X}^2 + \frac{614727}{40}\tilde{X}^3 - \frac{124389}{10}\tilde{X}^4 \\
& + \frac{355703}{80}\tilde{X}^5 - \frac{16568}{21}\tilde{X}^6 + \frac{7516}{105}\tilde{X}^7 - \frac{22}{7}\tilde{X}^8 + \frac{11}{210}\tilde{X}^9 \\
& + \tilde{S}^2 \left(-\frac{62391}{20} + \frac{614727}{20}\tilde{X} - \frac{1368279}{20}\tilde{X}^2 + \frac{4624139}{80}\tilde{X}^3 - \frac{157396}{7}\tilde{X}^4 \right. \\
& \quad \left. + \frac{30064}{7}\tilde{X}^5 - \frac{2717}{7}\tilde{X}^6 + \frac{2761}{210}\tilde{X}^7 \right) \\
& + \tilde{S}^4 \left(-\frac{124389}{10} + \frac{6046951}{160}\tilde{X} - \frac{248520}{7}\tilde{X}^2 + \frac{481024}{35}\tilde{X}^3 - \frac{15972}{7}\tilde{X}^4 \right. \\
& \quad \left. + \frac{18689}{140}\tilde{X}^5 \right) \\
& + \tilde{S}^6 \left(-\frac{70414}{21} + \frac{465992}{105}\tilde{X} - \frac{11792}{7}\tilde{X}^2 + \frac{19778}{105}\tilde{X}^3 \right) \\
& + \tilde{S}^8 \left(-\frac{682}{7} + \frac{7601}{210}\tilde{X} \right),
\end{aligned}$$

$$\tilde{S} \equiv \frac{X}{\sinh\left(\frac{X}{2}\right)}$$

Distortion of the spectral intensity

$$\Delta I \equiv \frac{X^3}{e^X - 1} \frac{\Delta n(X)}{n_0(X)}$$



Summary for thermal SZ effect

- Relativistic corrections are very important for high-temperature CG $T_e > 10\text{keV}$.
- Relativistic corrections are significant for large X region (sub-millimeter region)
For $X > 10$, factor 4 effect!
- Fokker-Planck expansion approximation is OK for $T_e < 15\text{keV}$.

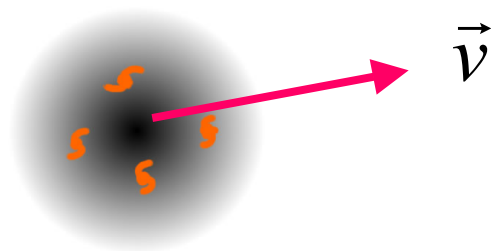
(b) Kinematical SZ Effect

Peculiar Motion of Clusters of Galaxies

Nozawa, Itoh & Kohyama ApJ, 508, 17 (1998)

Lorentz boosting of the original Boltzman equation

$$\left\{ \begin{array}{l} f(E) = f_{CG}(E_{CG}) , \quad f(E') = f_{CG}(E'_{CG}) \\ E_{CG} = \frac{E - \vec{\beta} \cdot \vec{p}}{\sqrt{1 - \beta^2}} , \quad E'_{CG} = \frac{E' - \vec{\beta} \cdot \vec{p}'}{\sqrt{1 - \beta^2}} , \quad \vec{\beta} = \frac{\vec{v}}{c} \end{array} \right.$$



Peculiar velocity of CG
relative to CMB.



Distortion of the CMB spectrum

$$\begin{aligned}
\frac{\Delta n(X)}{n_0(X)} &= \frac{\tau X e^X}{e^X - 1} \theta_e \left[Y_0 + \theta_e Y_1 + \theta_e^2 Y_2 + \theta_e^3 Y_3 + \theta_e^4 Y_4 \right] \\
&+ \frac{\tau X e^X}{e^X - 1} \beta^2 \left[\frac{1}{3} Y_0 + \theta_e \left(\frac{5}{6} Y_0 + \frac{2}{3} Y_1 \right) \right] \\
&+ \frac{\tau X e^X}{e^X - 1} \beta P_1(\hat{\beta}_z) \left[1 + \theta_e C_1 + \theta_e^2 C_2 \right] \\
&+ \frac{\tau X e^X}{e^X - 1} \beta^2 P_2(\hat{\beta}_z) \left[D_0 + \theta_e D_1 \right],
\end{aligned}$$

$$\hat{\beta}_z \equiv \frac{\beta_z}{\beta} = \cos \theta_\gamma,$$

$$P_1(\hat{\beta}_z) = \hat{\beta}_z,$$

$$P_2(\hat{\beta}_z) = \frac{1}{2} (3\hat{\beta}_z^2 - 1),$$

Kinematical SZ effect

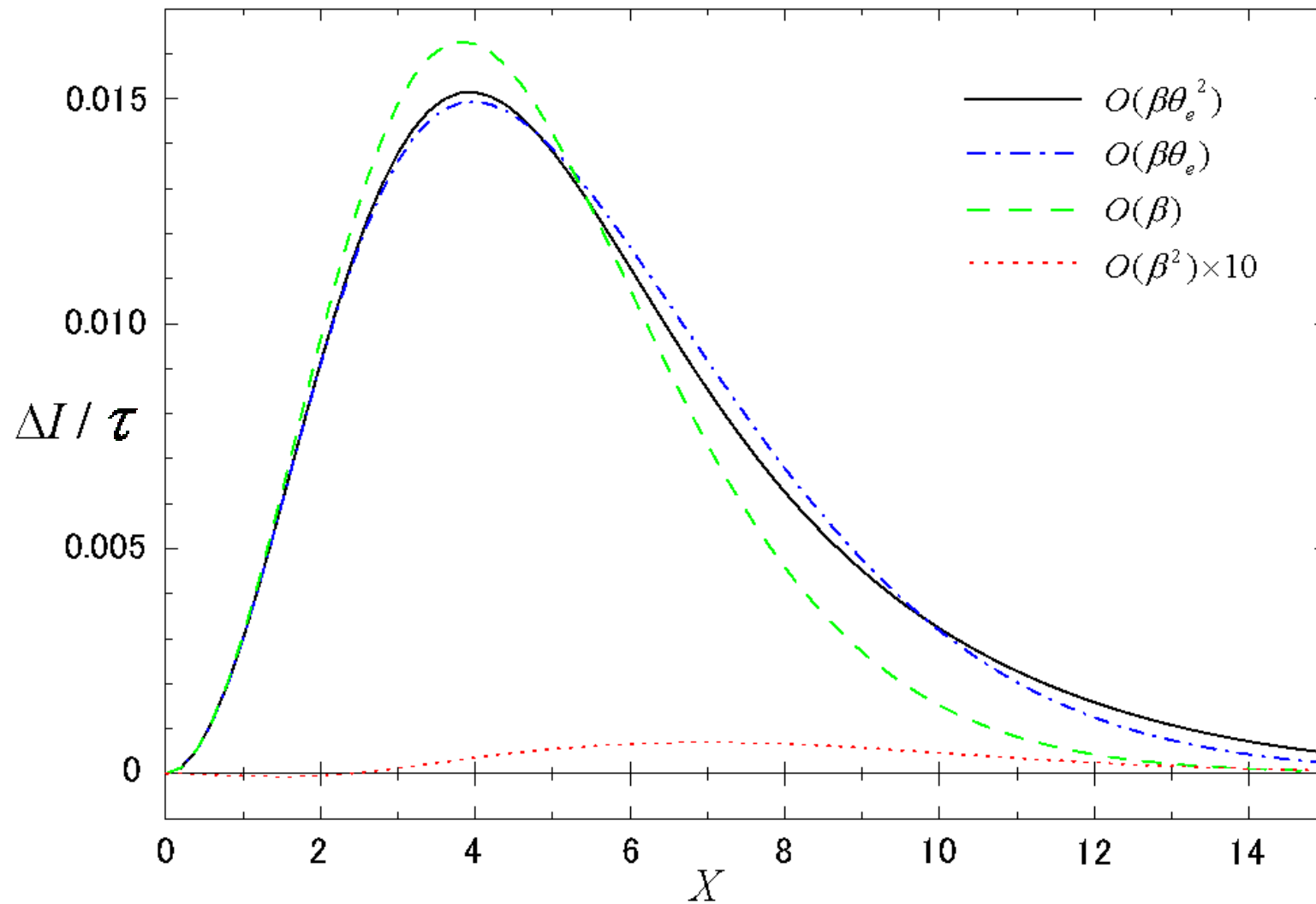
Relativistic corrections

$$\begin{aligned}
 C_1 &= 10 - \frac{47}{5}\tilde{X} + \frac{7}{5}\tilde{X}^2 + \frac{7}{10}\tilde{S}^2, \\
 C_2 &= 25 - \frac{1117}{10}\tilde{X} + \frac{847}{10}\tilde{X}^2 - \frac{183}{10}\tilde{X}^3 + \frac{11}{10}\tilde{X}^4 \\
 &\quad + \tilde{S}^2 \left(\frac{847}{20} - \frac{183}{5}\tilde{X} + \frac{121}{20}\tilde{X}^2 \right) + \frac{11}{10}\tilde{S}^4, \\
 D_0 &= -\frac{2}{3} + \frac{11}{30}\tilde{X}, \\
 D_1 &= -4 + 12\tilde{X} - 6\tilde{X}^2 + \frac{19}{30}\tilde{X}^3 + \tilde{S}^2 \left(-3 + \frac{19}{15}\tilde{X} \right).
 \end{aligned}$$

Distortion of the spectral intensity

$$\Delta I \equiv \frac{X^3}{e^X - 1} \frac{\Delta n(X)}{n_0(X)}$$

$k_B T_e = 10$ keV, $v = 1,000$ km/s



Summary for kinematical SZ effect (CG)

- Relativistic corrections are very important for high-temperature CG $T_e > 10\text{keV}$.

For $k_B T_e = 10\text{keV}$, $v = 1,000\text{km/s}$

$$O(\beta\theta_e) = -8.2\%$$

$$O(\beta\theta_e^2) = +1.3\%$$

Motion of the observer (Solar System)

Nozawa, Itoh & Kohyama submitted to A&A (2005)

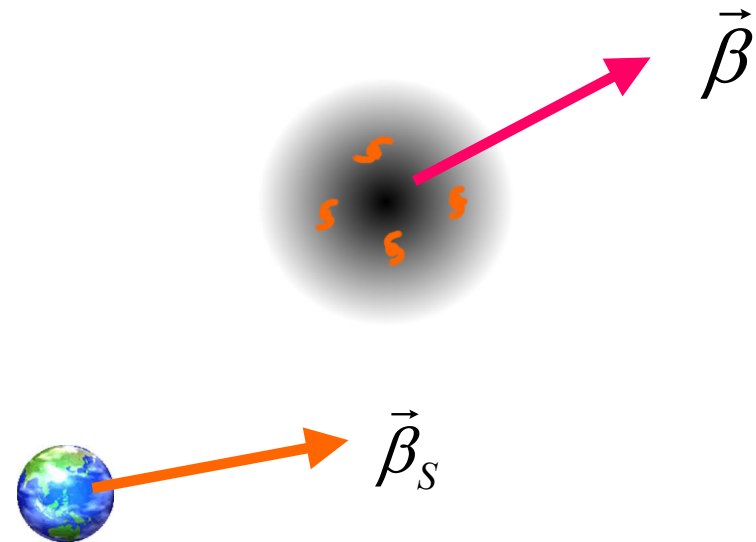
$\vec{\beta}_S$: velocity of the solar system respect to CMB

\vec{k}_S : photon vector in the solar system

$$n(\omega) = n(\omega_S),$$

$$\omega = \gamma_S (\omega_S + \vec{\beta}_S \cdot \vec{k}_S),$$

$$\gamma_S \equiv \frac{1}{\sqrt{1 - \beta_S^2}},$$



Planck distribution: dipolar distortion

$$n_S^0(X_S, \mu_S) = \frac{1}{\exp\{\gamma_S X_S (1 + \beta_S \mu_S)\} - 1},$$

$$X_S \equiv \frac{\omega_S}{k_B T_0},$$

$$\mu_S \equiv \hat{\beta}_S \cdot \hat{k}_S = \cos \theta_S,$$

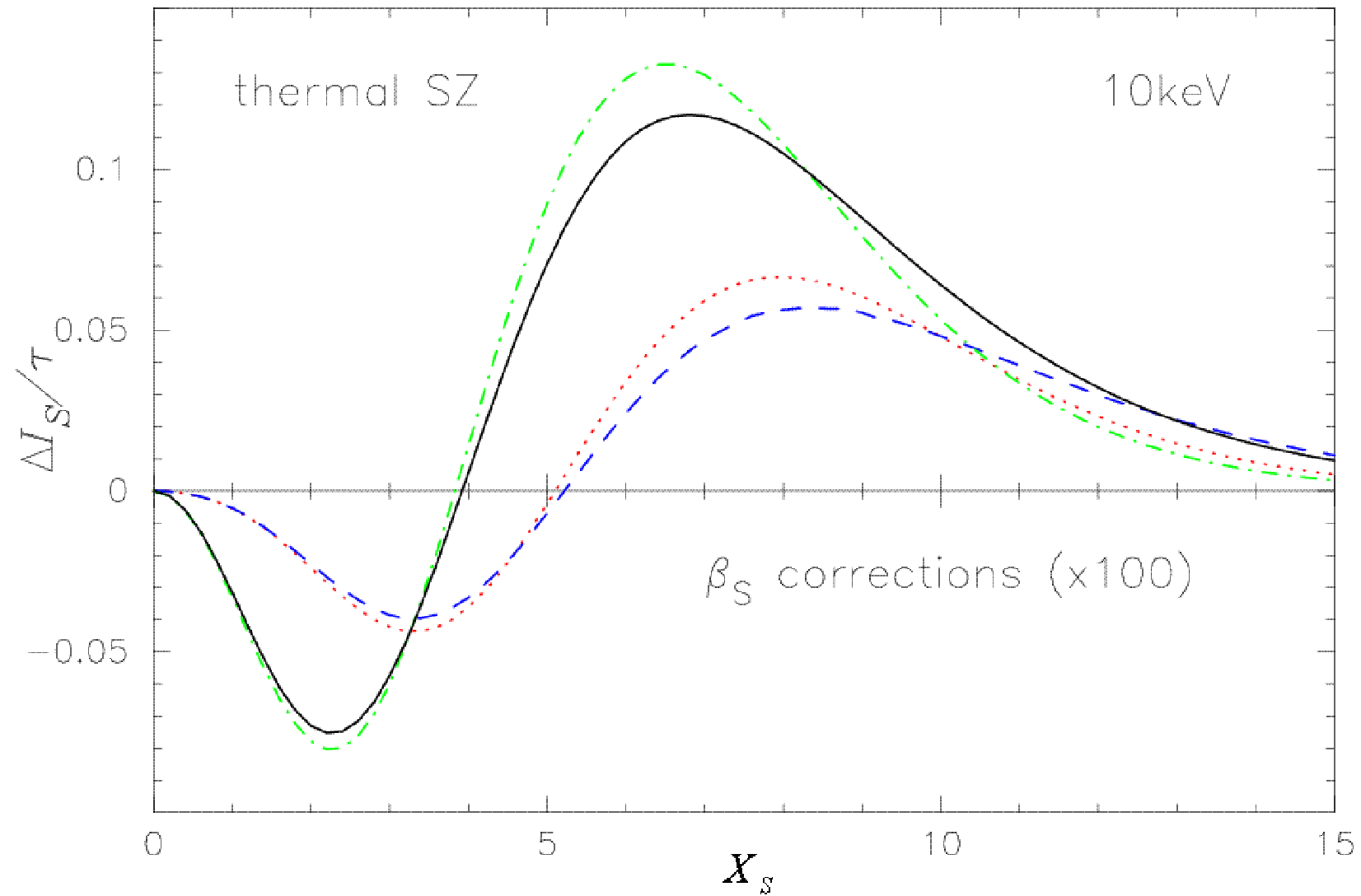
$$X = \gamma_S X_S (1 + \beta_S \mu_S)$$

velocity of CG in CMB frame is also transformed

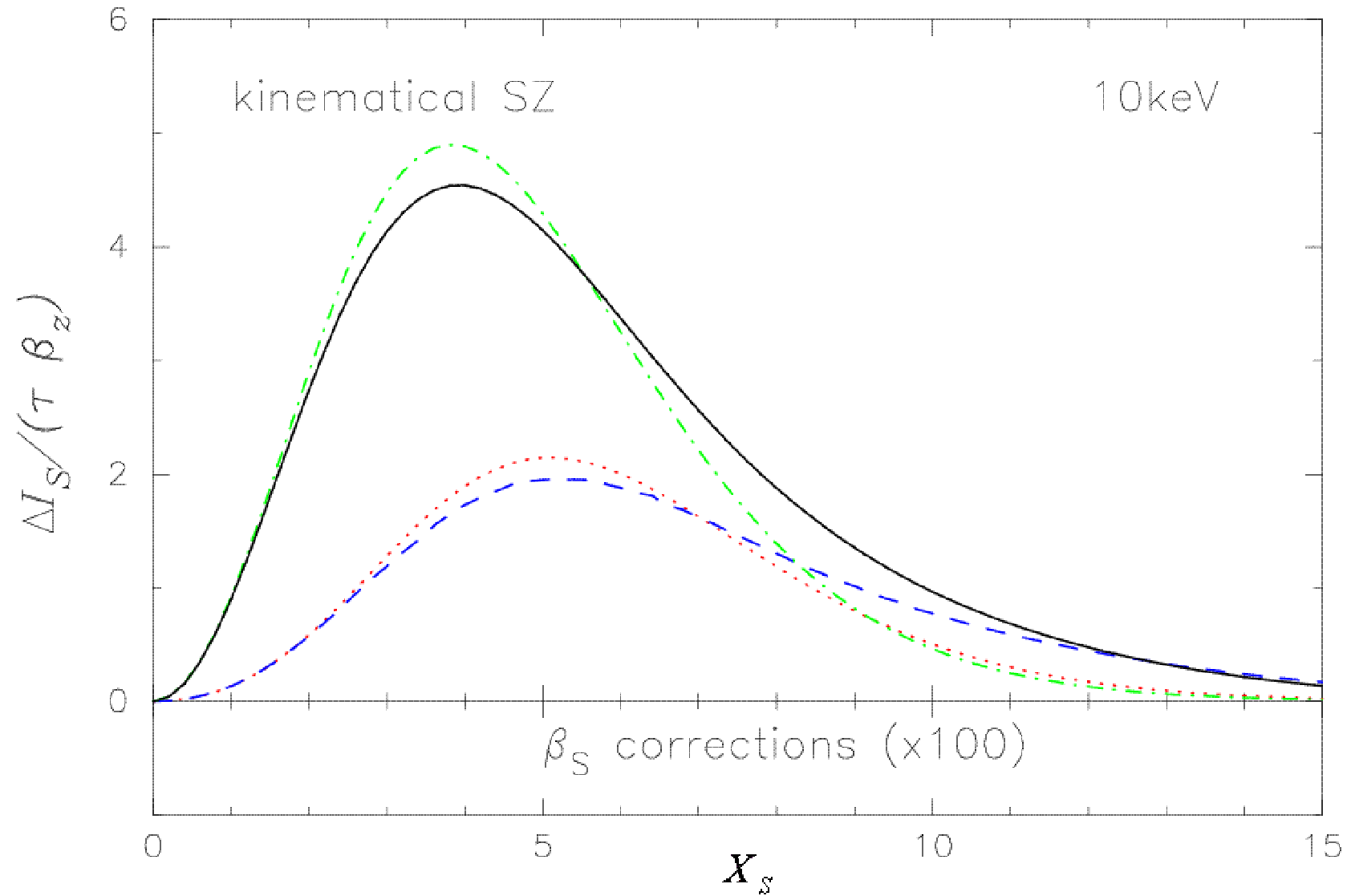
$$\vec{\beta} = \left(\frac{\vec{\beta}' \cdot \hat{\beta}_S + \beta_S}{1 + \vec{\beta}' \cdot \vec{\beta}_S} \right) \hat{\beta}_S + \frac{\vec{\beta}' - (\vec{\beta}' \cdot \hat{\beta}_S) \hat{\beta}_S}{\gamma_S (1 + \vec{\beta}' \cdot \vec{\beta}_S)}$$

$$\beta_z = \vec{\beta} \cdot \hat{k}$$

Thermal SZ (Effect of observer's motion)



Kinematical SZ (Effect of observer's motion)



Summary for kinematical SZ effect (Motion of the observer)

$$v_s = 372 \text{ km/s}$$

$$\beta_s = 1.24 \times 10^{-3}$$

For thermal SZ

β_s corrections are small

0.5% to 1% of the thermal SZ at $k_B T_e = 10 \text{ keV}$

For kinematical SZ

β_s corrections are small

0.5% to 1% of the thermal SZ at $k_B T_e = 10 \text{ keV}$

(c) Polarization SZ Effect

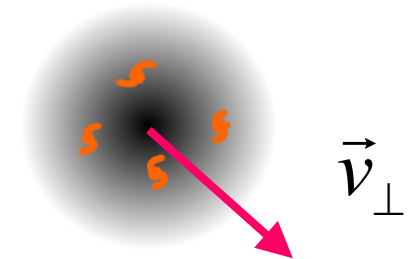
Itoh, Nozawa & Kohyama ApJ, 533, 588

(2000)

Polarized photons with Lorentz boosted Boltzman equation

$$\frac{\partial n(\omega)}{\partial t} = -2 \int \frac{d^3 p}{(2\pi)^3} d^3 p' d^3 k' W [1 + \vec{\eta} \cdot \vec{\zeta}]$$

$$\{n(\omega)[1 + n(\omega')]f(E) - n(\omega')[1 + n(\omega)]f(E')\},$$

 W : unpolarized crosssection $\vec{\eta} = (\eta_1, \eta_2, \eta_3)$: Stokes parameter $\vec{\zeta} = (\zeta_1, \zeta_2, \zeta_3)$: polarization crosssection

Magnitude of the polarization of CMB

$$P = \frac{\tau \theta_e X e^X}{e^X - 1} \beta_{\perp}^2 \left(F_0 + \theta_e F_1 + \theta_e^2 F_2 \right),$$

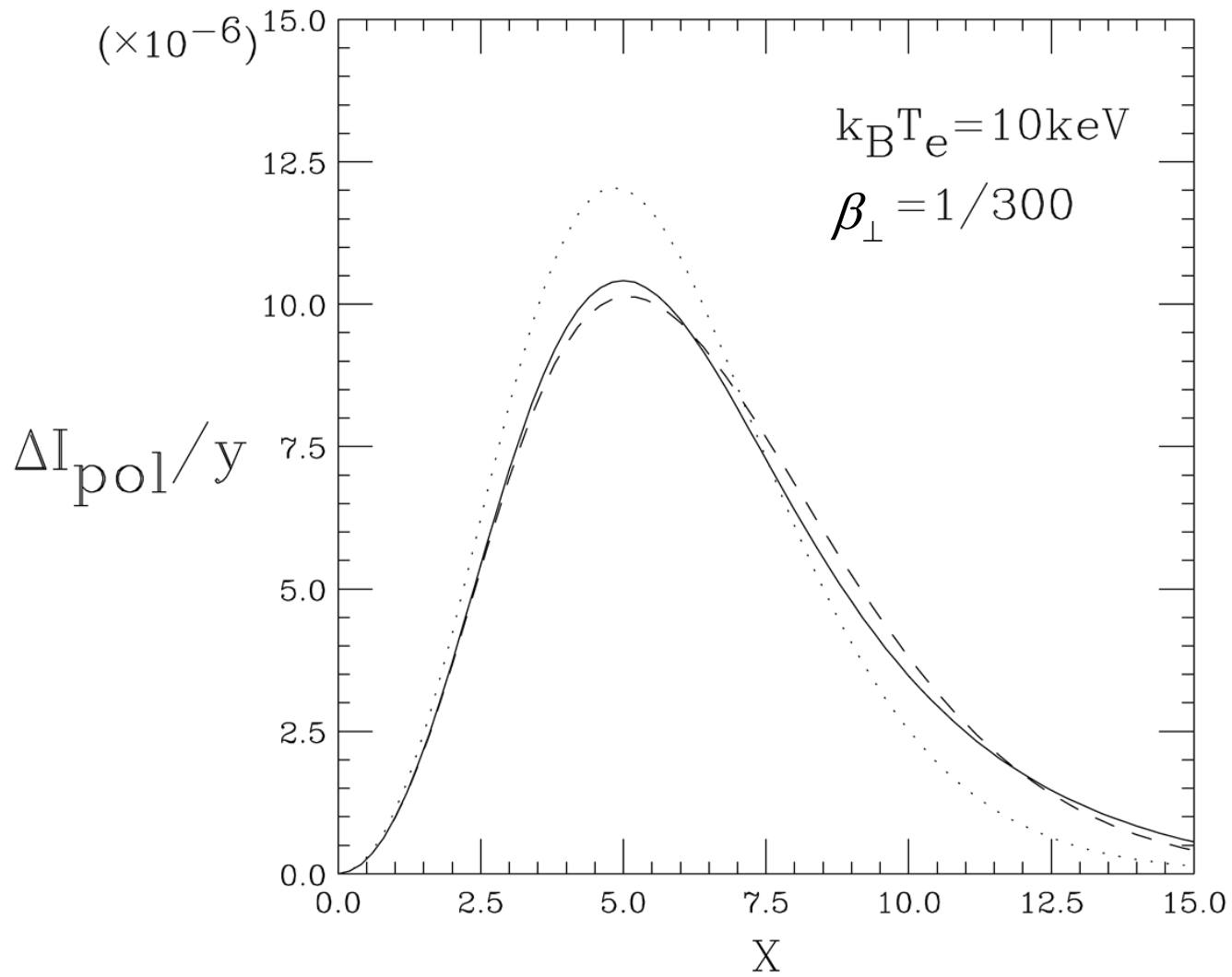
$$\tau \equiv \sigma_T \int dl N_e, \quad \theta_e = \frac{k_B T_e}{m_e c^2}$$

$$F_0 = \frac{1}{20} \tilde{X},$$

$$\left\{ \begin{aligned} F_1 &= \frac{3}{10} \tilde{X} - \frac{2}{5} \left(\tilde{X}^2 + \frac{1}{2} \tilde{S}^2 \right) + \frac{1}{20} \left(\tilde{X}^3 + 2 \tilde{X} \tilde{S}^2 \right), \\ F_2 &= \frac{3}{4} \tilde{X} - \frac{21}{5} \left(\tilde{X}^2 + \frac{1}{2} \tilde{S}^2 \right) + \frac{867}{280} \left(\tilde{X}^3 + 2 \tilde{X} \tilde{S}^2 \right) \\ &\quad - \frac{4}{7} \left(\tilde{X}^4 + \frac{11}{2} \tilde{X}^2 \tilde{S}^2 + \tilde{S}^4 \right) + \frac{1}{35} \left(\tilde{X}^5 + 13 \tilde{X}^3 \tilde{S}^2 + \frac{17}{2} \tilde{X} \tilde{S}^4 \right) \end{aligned} \right.$$

Relativistic corrections

Distortion of the spectral intensity



Summary for Polarization SZ effect

- Relativistic corrections are very important for high-temperature CG $T_e > 10\text{keV}$.
- Distortion of the spectral intensity is extremely small.

For $k_B T_e = 10\text{keV}$, $v = 1,000\text{km/s}$

$$\Delta I_{pol} / y = 1 \times 10^{-5}$$

(d) Double Scattering Effect

Assumption: initial photon distribution is isotropic

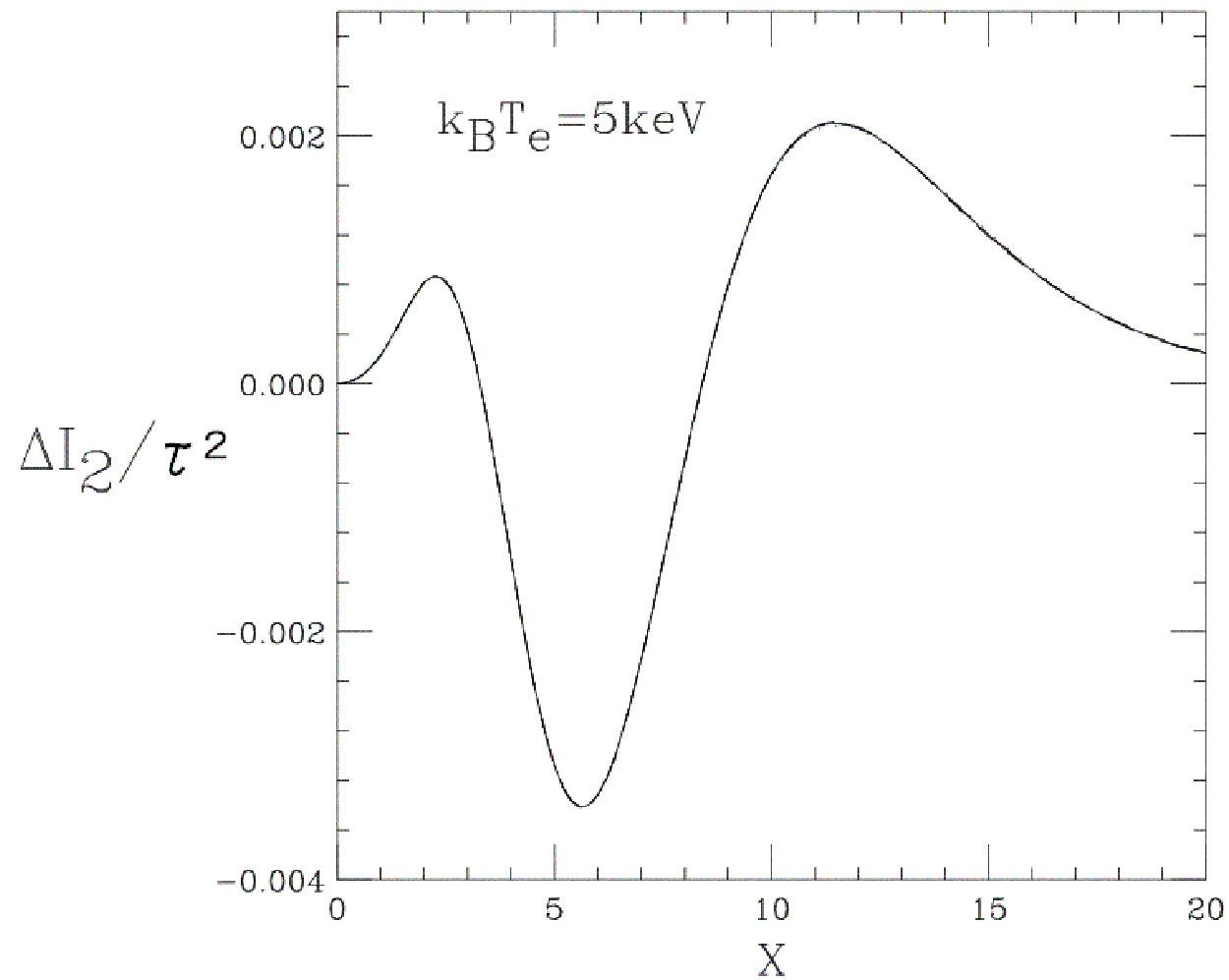
$$\begin{aligned}
 n(X) &= n_1(X) \equiv n_0(X) + \Delta n(X) \\
 &= n_0(X) \left[1 + \frac{\Delta n(X)}{n_0(X)} \right] \quad \text{—————} \quad (1)
 \end{aligned}$$

$$\frac{\Delta n(X)}{n_0(X)} = \frac{\tau \theta_e X e^X}{e^X - 1} [Y_0 + \theta_e Y_1 + \theta_e^2 Y_2 + \theta_e^3 Y_3 + \theta_e^4 Y_4]$$

Inserting equation (1) into Boltzman equation:

$$\Delta I = \Delta I_1 + \Delta I_2$$

$$\Delta I_2 = \frac{1}{2} \frac{\tau^2 \theta_e^2 X^4 e^X}{(e^X - 1)^2} (Z_0 + \theta_e Z_1 + \theta_e^2 Z_2 + \dots)$$



$\Delta I_2 / \Delta I_1 \approx -0.2$ per cent for $T_e = 15 \text{ keV}$

Double scattering effect is safely neglected.

(3) Summary

- We have presented a relativistically covariant formalism for the SZ effects.
- With the formalism one can treat thermal SZ effect, kinematical SZ effect (CG and observer), polarization SZ effect, and multiple scattering effect in an unified manner.
- Relativistic corrections are very important for high-temperature CG $T_e > 10\text{keV}$ for all SZ effects.